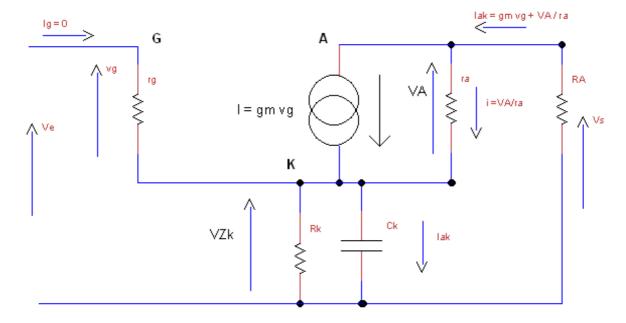
Common cathode triode amplifier Low cut-off frequency

common cathode amp stage with partially bypassed cathode



With:

- ra = internal resistance of the triode (cf Datasheet)
- VA = anode cathode voltage
- gm = transconductance (cf Datasheet)
- Zk = complex impedance on the cathode (Rk // Ck with Ck being a NOT perfect capacitance)
- RA load impedance on the anode
- rg = grid resistance held as infinite

The gain of the stage is:

$$Av = \frac{V_S}{Ve} = \frac{-\mu RA}{RA + ra + (\mu + 1)Zk}$$
(1)

This formula helps computing precisely the low cut-off frequency induced by the bypass capacitor Ck

In the bandwidth Zk = 0 and the gain is $Av = \frac{-\mu RA}{RA + ra}$

In fact
$$\mathbf{Z}\mathbf{k} = \mathbf{R}\mathbf{k} // \mathbf{C}\mathbf{k} = \frac{Rk}{1 + jCkRk} \omega$$

We see that for $\omega = \infty$ Zk = 0 and the gain is maximum. The cut-off frequency is confirmed as a low frequency and the problem becomes how low the value of ω can be before the gain reaches the value of $\mathbf{Av} - \mathbf{3} \, \mathbf{dB}$?

In fact for a same value of Ve (Vin) we will reach the low limit of the bandwidth when Vs (Vout) will be diminished by 3 dB, therefore at the cut-off $Vsc=\frac{Vs}{\sqrt{2}}$.

For this it is sufficient that the magnitude of the denominator of (1) equals $\sqrt{2}(RA+ra)$.

At the cut-off frequency ω_0 we have :

$$\sqrt{2}(RA + ra) = \begin{vmatrix} RA + ra + (\mu + 1) \frac{Rk}{1 + jCkRk} & \omega_0 \end{vmatrix}$$

$$\sqrt{2} = \frac{\begin{vmatrix} RA + ra + (\mu + 1) \frac{Rk}{1 + jCkRk} & \omega_0 \end{vmatrix}}{\begin{vmatrix} RA + ra \end{vmatrix}}$$

multiplied high and low by 1 + jCkRkw0 it yields:

$$\sqrt{2} = \frac{\left| (RA + ra)(1 + jCkRk \omega_0) + (\mu + 1)Rk \right|}{\left| (Ra + ra)(1 + jCkRk \omega_0) \right|}$$

$$\sqrt{2} = \frac{\left| (RA + ra) + (\mu + 1)Rk + j(RA + ra)RkCk \omega_0 \right|}{\left| RA + ra + j(RA + ra)RkCk \omega_0 \right|}$$

Hence simplified by Ra+ra:

$$\sqrt{2} = \frac{\left| 1 + \frac{(\mu + 1)Rk}{Ra + ra} + jRkCk \quad \omega_{0} \right|}{\left| 1 + jRkCk \quad \omega_{0} \right|}$$

Dividing the magnitudes yields:

$$\sqrt{2} = \frac{\sqrt{(1 + \frac{(\mu + 1)Rk}{Ra + ra})^2 + Rk^2Ck^2\omega_0^2}}{\sqrt{1 + Rk^2Ck^2\omega_0^2}}$$

Squaring it all:

$$2 + 2Rk^{2}Ck^{2}\omega_{0}^{2} = (1 + \frac{(\mu + 1)Rk}{Ra + ra})^{2} + Rk^{2}Ck^{2}\omega_{0}^{2}$$

$$2 + Rk^{2}Ck^{2}\omega_{0}^{2} = (1 + \frac{(\mu + 1)Rk}{Ra + ra})^{2}$$

$$\omega_{0}^{2} = \frac{(1 + \frac{(\mu + 1)Rk}{Ra + ra})^{2} - 2}{Rk^{2}Ck^{2}}$$

With $\omega_0 = 2 \pi$ f (Hz)

$$f_{(Hz)} = \frac{1}{2 \pi RkCk} \sqrt{\left[1 + \frac{(\mu+1)Rk}{RA + ra}\right]^2 - 2}$$

Of course if f is given, the value of the capacitance Ck can be found with the same formula, just reversing both.